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## A model of mitochondrial $\text{Ca}^{2+}$ -induced $\text{Ca}^{2+}$ release simulating the $\text{Ca}^{2+}$ oscillations and spikes generated by mitochondria

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### Abstract

Recent evidence underlines a key role of mitochondrial  $\text{Ca}^{2+}$  fluxes in cell  $\text{Ca}^{2+}$  signalling. We present here a kinetic model simulating the  $\text{Ca}^{2+}$  fluxes generated by mitochondria during mitochondrial  $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release (mCICR) resulting from the operation of the permeability transition pore (PTP). Our model connects the  $\text{Ca}^{2+}$  fluxes through the ruthenium red-sensitive  $\text{Ca}^{2+}$  uniporter, the respiration-dependent and passive  $\text{H}^+$  fluxes, the rate of oxygen consumption, the movements of weak acids across the mitochondrial membrane, the electrical transmembrane potential ( $\Delta\Psi$ ), and operation of the PTP. We find that two factors are crucial to account for the various mCICR profiles that can be observed experimentally: (i) the dependence of PTP opening and closure on matrix pH ( $\text{pH}_i$ ), and (ii) the relative inhibition of the respiratory rate consecutive to PTP opening. The resulting model can simulate irreversible  $\text{Ca}^{2+}$  efflux from mitochondria, as well as the genesis of damped or sustained  $\text{Ca}^{2+}$  oscillations, and of single  $\text{Ca}^{2+}$  spikes. The model also simulates the main features of mCICR, i.e. the threshold-dependence of mCICR triggering, and the all-or-nothing nature of mCICR operation. Our model should appear useful to further mathematically address the consequences of mCICR on the spatiotemporal organisation of  $\text{Ca}^{2+}$  signals, as a ‘plug-in’ module for the existing models of cell  $\text{Ca}^{2+}$  signalling. © 1998 Elsevier Science B.V. All rights reserved

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### 1. Introduction

In many cells, the information provided by external stimuli, either electrical, hormonal, photochemical, or

mechanical, is conveyed within the cytoplasm by  $\text{Ca}^{2+}$  signals that take the form of  $\text{Ca}^{2+}$  oscillations and waves [1]. The information is encoded by the spatiotemporal organisation of these signals, and the latter determines the subsequent cellular responses [2]. The main mechanism underlying the organisation of the cytosolic  $\text{Ca}^{2+}$  signals is autocatalytic and termed ‘ $\text{Ca}^{2+}$ -induced  $\text{Ca}^{2+}$  release’ (CICR) [1,3,4]. The

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CICR machinery is usually identified with the sole endoplasmic reticulum (ER), and thought to rely principally on  $\text{Ca}^{2+}$ -activable  $\text{Ca}^{2+}$  channels of the ER membrane, i.e. the inositol 1,4,5-trisphosphate ( $\text{IP}_3$ ) and ryanodine receptors ( $\text{IP}_3\text{R}$  and  $\text{RyR}$ , respectively) [1,3,4]. However, when  $\text{Ca}^{2+}$  is released from the ER at locally high concentration, the neighbouring mitochondria take up some of it as shown by aequorin addressed to mitochondria [5]. This results in a modification of the spatiotemporal organisation of the  $\text{Ca}^{2+}$  signals [6–8]. Two types of mechanisms can be proposed to account for such mitochondrial effect: (i) a regulation of  $\text{IP}_3\text{R}$  activity by local mitochondrial  $\text{Ca}^{2+}$  buffering, modulating the ER CICR [6], and (ii) a  $\text{Ca}^{2+}$ -stimulated operation of mitochondrial  $\text{Ca}^{2+}$  efflux pathways contributing to the cytosolic  $\text{Ca}^{2+}$  raise [9,11]. Regarding this last interpretation, one pathway of mitochondrial  $\text{Ca}^{2+}$  efflux is of particular interest, since it also obeys autocatalytic rules. This pathway, whose operation we termed mCICR [9,10], depends on the transient opening of the PTP by mitochondrial  $\text{Ca}^{2+}$  uptake, and generates  $\text{Ca}^{2+}$  spikes [9,10] and oscillations [12,13] in vitro. It should be noticed that in all these experimental observations, the external calcium concentration necessary to trigger oscillations or a spike on isolated mitochondria or mitochondria inside permeabilised cells are rather high, between 60 and 120  $\mu\text{M}$  and that, in these conditions, mitochondria are able to accumulate a large amount of the added calcium. Damped oscillations are induced by 100  $\mu\text{M}$  external calcium on permeabilised Ehrlich ascites tumour cells [12]; the sensitivity of these oscillations to cyclosporin A indicates the participation of PTP in this process. Sustained oscillations are observed with digitonin-treated *Tetrahymena pyriformis* in low  $\text{Po}_2$  conditions, i.e. at low respiratory rate [13]. In living cells, mCICR is triggered by  $\text{IP}_3$ -dependent  $\text{Ca}^{2+}$  release from the ER, resulting in the amplification of the cytosolic  $\text{Ca}^{2+}$  signals [9,10], indicating that mitochondria inside cells are either more sensitive to calcium or that high local calcium concentrations can be sensed by mitochondria in the vicinity of ER as demonstrated in [5], or both.

We present here a theoretical analysis of mCICR providing direct information on the mitochondrial determinants of mCICR. We based our mathematical model on the following observations: (i) extra-

mitochondrial  $\text{Ca}^{2+}$  is taken-up electrophoretically by mitochondria through the ruthenium red-sensitive  $\text{Ca}^{2+}$  uniporter [14]; (ii) mitochondrial  $\text{Ca}^{2+}$  uptake causes a transitory drop in  $\Delta\Psi$  that results in respiration activation [14]; (iii) activation of respiration, which corresponds to an increased  $\text{H}^+$  extrusion, results in the alcalinization of the mitochondrial matrix [14]; (iv) matrix alcalinization activates the PTP [15], and thus triggers various ion fluxes – among those,  $\text{Ca}^{2+}$  efflux – according to their electrochemical potential; (v) PTP activation also causes a relative respiration inhibition [12] that contributes to matrix acidification; (vi) matrix acidification inactivates the PTP [15].

The resulting model, although it is a very simplified theoretical model, simulates all the known modes and features of mCICR and should thus appear useful to perform mathematical studies of the mitochondrial contribution to cell  $\text{Ca}^{2+}$  signals as a ‘plug-in’ module for the already existing models of  $\text{Ca}^{2+}$  signalling [16–19].

## 2. The model

The model consists of a system of differential equations describing the following variables: outside  $\text{Ca}^{2+}$  concentration ( $\text{Ca}$ ),  $\Delta\Psi$  ( $\Delta\Psi = \Psi_{\text{out}} - \Psi_{\text{in}}$ ), inside  $\text{H}^+$  concentration ( $H$ ), the amount of opened PTP channels responsible for unselective mitochondrial membrane permeability ( $T^{\text{op}}$ ), matrix concentration of a respiratory chain activator ( $A$ ).

The model accounts for the following fluxes across the inner mitochondrial membrane:  $J_{\text{ca}}$ ,  $\text{Ca}^{2+}$  flux through the ruthenium red-sensitive mitochondrial  $\text{Ca}^{2+}$  uniporter;  $J_{\text{an}}$ , anions flux;  $J_{\text{AH}}$ , flux of the protonated acids;  $J_{\text{H}}$ ,  $\text{H}^+$  leak;  $J_{\text{r}}$ ,  $\text{H}^+$  transport connected with respiration;  $V_{\text{A}}$ , rate of production of a respiration activator within the matrix;  $J_{\text{A}}$ , efflux of this activator.

We have not included in the model the calcium transport through the activity of the  $2\text{Na}^+/\text{Ca}^{2+}$  nor through the  $\text{Ca}^{2+}2\text{H}^+$  exchangers, because experimentally, spikes have been shown to be generated independently of these exchangers (no  $\text{Na}^+$  in the medium) [9,10]. However, it should be noticed that experimental calcium oscillations have been observed in mitochondria in the presence of 100 mM Na [12] which

could indicate a role played by the  $2\text{Na}^+/\text{Ca}^{2+}$  exchanger. Thus our model will appear as a minimal model able to generate calcium oscillations and spikes in mitochondria; for the same reason the ATPase activity is not included.

We use the Goldman equation to describe the various ion fluxes through the inner mitochondrial membrane:

$$J = J^{\text{out}} - J^{\text{in}} = P\Delta\Psi \frac{C_{\text{in}} \exp\left(-zF \frac{\Delta\Psi}{RT}\right) - C_{\text{out}}}{1 - \exp\left(-zF \frac{\Delta\Psi}{RT}\right)} \quad (1)$$

where  $J^{\text{in}}$  denotes the ion flux from outside to inside, and  $J^{\text{out}}$  the opposite flux of the same ion.  $C_{\text{in}}$  and  $C_{\text{out}}$  are the ion concentrations inside and outside.  $P = p/(RT/zF)$ , where  $p$  is the permeability coefficient of the considered ion,  $z$  its charge,  $F$  is Faraday's number,  $R$  is the gas constant, and  $T$  the absolute temperature. In the following text we will use the term 'permeability coefficient' to qualify  $P$ .

### 2.1. $\text{Ca}^{2+}$ fluxes through the $\text{Ca}^{2+}$ uniporter

The  $\Delta\Psi$  dependence of the  $\text{Ca}^{2+}$  fluxes through the uniporter fits a form consistent with the Goldman constant field equation [20] and therefore can be described by Eq. (1). However, the  $\text{Ca}^{2+}$  fluxes mediated by the uniporter are not only dependent on transmembrane potential  $\Delta\Psi$ , but also on the external  $\text{Ca}^{2+}$  concentration. This is the reason why we included in the Goldman equation the term  $(\text{Ca} - C_0)$  instead of  $\text{Ca}$  only, where  $\text{Ca}$  represents extramitochondrial  $\text{Ca}^{2+}$  concentration and  $C_0 = 0.5 \mu\text{M}$ , which is the basal  $\text{Ca}^{2+}$  concentration above which the uniporter insures a net influx of  $\text{Ca}^{2+}$  [20,22].

$$J_{\text{Ca}} = P_{\text{Ca}}\Delta\Psi \frac{C_{\text{in}} \exp\left(-2F \frac{\Delta\Psi}{RT}\right) - (\text{Ca} - C_0)}{1 - \exp\left(-2F \frac{\Delta\Psi}{RT}\right)} \quad (2)$$

mCICR is triggered in vitro by extramitochondrial  $\text{Ca}^{2+}$  concentrations in the range of 60–100  $\mu\text{M}$ . In these conditions, the intra-mitochondrial  $\text{Ca}^{2+}$  exchangeable with outside  $\text{Ca}^{2+}$ , determined by null-point titration experiments, is in the same order of

magnitude than the outside one [21]. Moreover, because of the high intra-mitochondrial  $\text{Ca}^{2+}$  buffering power [14], we consider the internal  $\text{Ca}^{2+}$  concentration as a constant equal to 100  $\mu\text{M}$  ( $\text{Ca}_{\text{in}} = 100 \mu\text{M}$ ). It should be noted that the actual free calcium concentration is much lower and can be slightly variable.

$P_{\text{Ca}}$  can be determined from the measured value of the maximal  $\text{Ca}^{2+}$  flux through the uniporter. The maximal value of 1000  $\text{nmol min}^{-1} \text{mg}^{-1}$  [14,23] can be reproduced by Eq. (2) at  $\Delta\Psi = 150 \text{ mV}$  by taking  $P_{\text{Ca}} = 0.1 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ . However, the experimental value for the maximal rate of  $\text{Ca}^{2+}$  uptake is obtained at a lower  $\Delta\Psi$ , because it is measured in the presence of high outside  $\text{Ca}^{2+}$  concentration which collapses  $\Delta\Psi$ . Thus, the highest reported values of  $V_{\text{max}}$  for the  $\text{Ca}^{2+}$  uniporter rather reflect the maximal rate of respiration than the maximal rate of  $\text{Ca}^{2+}$  entry. Thus the real maximal  $\text{Ca}^{2+}$ -transporting capacity of uniporter can be considered as being higher [14]. For this reason, we set in the model  $P_{\text{Ca}} = 1 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ .

### 2.2. Opening and closure of the PTP

Regulation of the PTP is one of the key elements of our model. Three main factors have been shown to induce PTP opening: (i) increase in inside  $\text{Ca}^{2+}$  concentration; (ii) decrease in  $\Delta\Psi$ ; and (iii) increase in  $\text{pH}_i$  [15]. Conversely, the opposite changes of the same parameters promote PTP closure. mCICR can be induced by  $\text{Sr}^{2+}$ , which is not a direct agonist of PTP [22]. Thus, we make the two following assumptions: (i)  $\text{Ca}^{2+}$  itself is not the trigger for PTP opening during mCICR; and (ii) the amount of  $\text{Ca}^{2+}$  permanently present inside mitochondria is sufficient to promote PTP opening by  $\text{pH}_i$  or  $\Delta\Psi$ . Thus, according to our hypotheses, PTP activation during mCICR is indirectly triggered by  $\text{Ca}^{2+}$  uptake, but not by  $\text{Ca}^{2+}$  per se. Since  $\text{Ca}^{2+}$  uptake decreases  $\Delta\Psi$  and raises  $\text{pH}_i$ , these two factors are still candidate triggers of PTP opening during mCICR. However, the experiments show that during mCICR, PTP opening does not coincide with the minimum  $\Delta\Psi$  encountered during the mCICR cycle, and that PTP closure occurs when  $\Delta\Psi$  reaches its lowest value [9,10]. Thus, we suppose in our model that  $\Delta\Psi$  is not the trigger for PTP opening during mCICR. Regarding the role of  $\text{pH}_i$ , the max-

imum probability of PTP opening is observed at  $\text{pH}_i$  above 7.3, while the PTP remains closed at  $\text{pH}_i$  below 7.0 [15]. During mCICR, we observed experimentally that  $\text{pH}_i$  is maximal at the time of PTP opening, and minimal at the time of PTP closure [10]. Thus, we consider in our model that  $\text{pH}_i$  is the only key PTP regulator during mCICR. This assumption will appear below sufficient to account for the various mCICR profiles experimentally observed.

We thus utilise a function  $f$  which expresses the effect of  $\text{pH}_i$  on PTP opening:  $f = 1$  if  $\text{pH} > 7.5$ ,  $f = 0$  if  $\text{pH} < 7.0$  and  $f$  linearly increases from 0 to 1 when  $\text{pH}_i$  increases from 7.0 to 7.5. With this function, the rate of PTP opening ( $V_T^{\text{op}}$ ) is proportional to the amount of closed channels ( $1 - T^{\text{op}}$ ), and to the value of the function  $f$ . Conversely, the rate of PTP closure ( $V_T^{\text{cl}}$ ) is proportional to the value of  $(1 - f)$  and to the amount of opened channels ( $T^{\text{op}}$ ).

Thus, the rates are expressed by the equations:

$$V_T^{\text{op}} = k_T^{\text{op}} f (1 - T^{\text{op}}) \quad (3)$$

$$V_T^{\text{cl}} = k_T^{\text{cl}} (1 - f) T^{\text{op}} \quad (4)$$

The rate constants are supposed to be equal to  $1 \text{ min}^{-1}$ .

### 2.3. Electrogenic ion fluxes through the PTP

The  $\text{H}^+$  fluxes through the PTP are described using the general Eq. (1), taking also into account the PTP opening status ( $T^{\text{op}}$ ):

$$J_{\text{H}} = P_{\text{H}} T^{\text{op}} \Delta \Psi \frac{H \exp\left(-F \frac{\Delta \Psi}{RT}\right) - H_{\text{out}}}{1 - \exp\left(-F \frac{\Delta \Psi}{RT}\right)} \quad (5)$$

In experimental mCICR, outside  $\text{H}^+$  concentration ( $H_{\text{out}}$ ) is kept constant by pH buffers [9,12], thus we keep it constant in our model:  $H_{\text{out}} = 0.1 \text{ } \mu\text{M}$  ( $\text{pH} = 7.0$ ). Inside  $\text{H}^+$  concentration ( $H$ ) is variable. The value of  $P_{\text{H}}$ , the coefficient of permeability of the mitochondrial membrane for protons is unknown, but it can be assumed that the  $\text{H}^+$  fluxes during maximal PTP opening are of the same order as the  $\text{Ca}^{2+}$  fluxes. Thus,  $P_{\text{H}}$  can easily be estimated. Initial  $\text{Ca}^{2+}$  concentration pulse is set in our simulation to  $30 \text{ } \mu\text{M}$ , which is 300 times greater than outside  $\text{H}^+$  concentration.

Thus, with  $P_{\text{H}} = 300 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ , the two fluxes are of the same order.

Anion fluxes can be described with the same general Eq. (1), but the following considerations simplify their mathematical description:

Taking into consideration the equilibrium equation for acid dissociation:

$$\text{AH} \leftrightarrow [\text{A}^-] + [\text{H}^+], \quad K = \frac{[\text{A}^-][\text{H}^+]}{[\text{AH}]}$$

and the conservation equation:

$[\text{A}^-] + [\text{AH}] = \text{An}$ , where An is the total acid concentration, we get the following expression for  $[\text{A}^-]$ :

$$[\text{A}^-] = \frac{K \text{An}}{K + [\text{H}^+]}$$

At  $\text{pH} > 7$ ,  $K \gg [\text{H}^+]$ , and thus, we can consider that anion concentration is equal to the total concentration of all acids and their soluble salts,  $[\text{A}^-] = \text{An}$ .

Total acid concentration (An), is usually high compared to its changes during mCICR. This is due to the fact that anion variations outside mitochondria are damped by dilution (outside volume = 1000 times matrix volume), and changes in matrix anion concentrations are compensated by osmotic regulation.

Therefore, we will consider anion concentration, inside as well as outside, as a constant  $\text{An} = 10 \text{ mM}$ . In these conditions, Eq. (1) for anions simplifies as:

$$J_{\text{An}} = -P_{\text{An}} T^{\text{op}} \Delta \Psi \text{An} \quad (6)$$

Concerning  $P_{\text{An}}$ , the same reasoning as above leads to:  $P_{\text{An}} = 0.001 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ .

Regarding  $\text{Ca}^{2+}$ , we do not take into account the specific efflux of  $\text{Ca}^{2+}$  through the PTP because when  $\Delta \Psi$  collapses, the efflux of  $\text{Ca}^{2+}$  through the uniporter is only limited by proton influx and anion efflux through the PTP. We have experimentally shown that the dissipation of  $\Delta \Psi$  leads to  $\text{Ca}^{2+}$  efflux equally well through the uniporter or through the PTP, because it is not blocked by ruthenium-red or cyclosporin A added separately, but only by the presence of both [22]. For the sake of simplicity, we have only considered in the model the  $\text{Ca}^{2+}$  efflux through the uniporter. Thus Eq. (2) is sufficient to account for all  $\text{Ca}^{2+}$  movements through the mitochondrial mem-

brane, independently of the  $\text{Ca}^{2+}$ -conductive pathway provided by the PTP.

#### 2.4. $\text{H}^+$ ejection by the respiratory chain

The respiration rate, and the corresponding  $\text{H}^+$  efflux rate, is constant and maximal if the value of electrochemical transmembrane potential, expressed in electrical units,  $(\Delta\mu_{\text{H}^+} = \Delta\Psi + (RT/F)\Delta\text{pH}; \Delta\text{pH} = \text{pH}_i - \text{pH}_{\text{out}})$  does not exceed 160 mV. The maximal rate ( $J_{\text{max}}$ ) is approximately equal to 1000  $\text{nmol H}^+ \times \text{min}^{-1} \times \text{mg protein}^{-1}$ . If  $\Delta\mu_{\text{H}^+} > 160$  mV, we consider that the flux of protons linearly decreases together with the increase of  $\Delta\mu_{\text{H}^+}$ , and that the flux falls to zero when  $\Delta\mu_{\text{H}^+}$  reaches 200 mV.

Thus,

$$\text{If } \Delta\mu_{\text{H}^+} > 160\text{mV, } J'_r = J_{\text{max}} \left(1 - \frac{(\Delta\mu_{\text{H}^+} - 160)}{40}\right),$$

$$\text{otherwise } J'_r = J_{\text{max}} \quad (7)$$

It was shown that the respiration rate undergoes a relative inhibition consecutive to PTP opening, while  $\Delta\Psi$  is decreased [12]. The factor, which induces this inhibition is still unknown. We thus suppose here the existence of an activator of respiration, such that the respiration rate is proportional to its intra-mitochondrial concentration:

$$J_r = J'_r A \quad (8)$$

where  $J'_r$  is described by Eq. (7), and  $A$  is the amount of activator inside the mitochondrial matrix. We do not speculate here on the very nature of the activator, and use it as a phenomenological variable.

We further assume that the relative inhibition of respiration during PTP opening is due to the efflux of activator through the opened PTP. The efflux of activator ( $J_A$ ) is proportional to its relative amount inside mitochondria, and to the PTP opening state:

$$J_A = P_A [T^{\text{op}}][A] \quad (9)$$

where the permeability coefficient  $P_A$  is set to 1  $\text{min}^{-1}$ .

The production of activator ( $V_A$ ) is supposed to be dependent on the relative amount of activator inside mitochondria (between 0 and 1). The production is stopped when the activator concentration reaches the relative maximal level of 1 and it linearly increases when the concentration decreases:

$$V_A = k_A(1 - A) \quad (10)$$

where  $k_A$  is a constant set to 1  $\text{min}^{-1}$ .

To represent respiration in terms of  $\text{O}_2$  consumption, we fixed a finite initial oxygen pool of 500  $\text{nmol O}_2 \text{ mg protein}^{-1}$ , and assumed a rate of consumption equal to  $-J_r/4$ .

#### 2.5. $e - \Delta\Psi$

$\Delta\Psi$  changes in response to the described electrogenic ion fluxes correspond to:

$$\frac{d(\Delta\Psi)}{dt} = \frac{F}{C}(J_r + 2J_{\text{Ca}} + J_{\text{An}} + J_{\text{H}}) \quad (11)$$

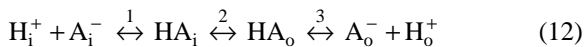
where  $F = 10^5 \text{ q/mol} = 10^{-4} \text{ q/nmol}$  (Faraday number) and  $C$  is the electric capacity of the membrane.

The measured specific capacity of biomembranes approaches 1  $\text{mF/cm}^2$  [24]. The mitochondrial surface is function of the osmotic pressure, and varies from 1  $\text{m}^2/\text{g}$  of protein to 30  $\text{m}^2/\text{g}$  of protein [25]. Experimentally, mCICR operation is favoured by hypo-osmotic conditions [9], due to the sensitisation of the PTP to  $\text{Ca}^{2+}$  [26]. We thus set here the value of  $S$  to 20  $\text{m}^2/\text{g} = 200 \text{ cm}^2/\text{mg}$ . Therefore  $C = 200 \text{ mF/mg of protein} = 2 \times 10^{-4} \text{ F/mg of protein}$ . So,  $F/C = 0.5 \text{ V nmol}^{-1} \text{ mg protein} = 500 \text{ mV nmol}^{-1} \text{ mg protein}$ .

#### 2.6. Electroneutral fluxes of weak acids

During mCICR, the fluxes of weak organic acids, i.e. principally succinate under the reference in vitro conditions [9,12], allow the recovery of  $\text{pH}_i$  after  $\text{H}^+$  respiratory efflux, due to the dissociation of the transported acids inside the matrix. In the particular case of succinate, the latter is transported through the mitochondrial membrane by a dicarboxylate translocator, in exchange for malate. Using the same translocator, malate goes back into the matrix in exchange for phosphate. Finally the operation of the phosphate transporter exchanges external phosphate for matrix  $\text{OH}^-$ . In the reference experimental conditions it appears that succinate transport is never a limiting step. Thus, we simplify here the mathematical description of the weak acid fluxes by assuming that succinate is membrane-permeant under its protonated form.

The electroneutral fluxes of permeant weak acids proceed according to the scheme:



Steps 1 and 3 correspond to the reversible dissociation of the acids, inside (i) and outside (o) mitochondria, step 2 corresponds to the diffusion through the membrane. The rate of diffusion ( $J$ ) is proportional to the protonated acid (HA) concentrations inside and outside:

$J = k[\text{AH}]$ , or taking into account equilibrium steps 1 and 3,  $J = k_1[\text{A}^-][\text{H}^+]$ , where  $k_1 = k/K$ .

But since  $[\text{A}^-] = \text{An}$ , as shown before,  $J = k_{\text{AH}}[\text{H}^+]$ , where  $k_{\text{AH}} = k_1\text{An}$ .

Thus, if inside and outside anion concentrations are equal, as well as the constants of diffusion through the membrane in both directions, the equation giving the net flux of AH is:

$$J_{\text{AH}} = J^{\text{out}} - J^{\text{in}} = k_{\text{AH}}(\text{H} - \text{H}_{\text{out}}) \quad (13)$$

where we set  $k_{\text{AH}}$  to  $100 \text{ min}^{-1} \text{ mg protein}^{-1}$ , and  $\text{H}_{\text{out}}$  to  $0.1 \mu\text{M}$  ( $\text{pH} = 7$ ).

Other fluxes such as  $\text{Ca}^{2+}/\text{H}^+$ ,  $\text{Ca}^{2+}/\text{Na}^+$ , or  $\text{Na}^+/\text{H}^+$  exchanges, are not taken into account in the model, because they are exceedingly slow in comparison with the fluxes encountered during mCICR (for instance, the rate of  $\text{Na}^+/\text{H}^+$  exchange is only  $0.5 \text{ nmol min}^{-1} \text{ mg protein}^{-1}$  [14]).

### 2.7. Differential equations of the model

The following system of differential equations describes the behaviour of the variables:

$$\begin{aligned} \frac{d(\text{Ca})}{dt} &= J_{\text{Ca}}/v_{\text{out}} \\ \frac{d(\Delta\Psi)}{dt} &= \frac{F}{C}(J_r + 2J_{\text{Ca}} + J_{\text{AN}} + J_{\text{H}}) \\ \frac{d(\text{H})}{dt} &= (-J_{\text{H}} - J_{\text{AH}} - J_r)/B/v_{\text{in}} \\ \frac{d(T^{\text{op}})}{dt} &= V_T^{\text{op}} - V_T^{\text{cl}} \\ \frac{d(\text{A})}{dt} &= V_{\text{A}} - J_{\text{A}} \end{aligned} \quad (14)$$

The ion fluxes are also dependent on the amount of mitochondria undergoing mCICR. On the basis of

experimental conditions [9,12] we set it to 1 mg of proteins. The total outside volume  $V_{\text{out}}$  is set to 1 ml, and the matrix volume  $V_{\text{in}}$  to  $2 \mu\text{l}$ . The  $\text{H}^+$  buffering capacity of the mitochondrial matrix,  $B$ , can be estimated by the following consideration: the uptake of  $\text{Ca}^{2+}$ , at an outside concentration of  $20\text{--}50 \mu\text{M}$ , leads to an equivalent efflux of  $\text{H}^+$ , which corresponds to a matrix  $\text{H}^+$  concentration decrease inferior to  $0.1 \mu\text{M}$ . Taking into account that the matrix volume is about three orders of magnitude lower than the outside one, we set the buffer capacity to  $B = 300\,000$ .

The numeric solution of the system of differential equations was performed using the software 'Pspice', aimed at solving rigid systems of differential equation, or with a home-made software using the method of backward integration with variable steps. The two methods provided identical solutions.

## 3. Results and discussion

The model described above, and schematically represented in Fig. 1, takes into account two main experimental results: (i) PTP opening and closure are induced by changes in  $\text{pH}_i$  [15]; and (ii) PTP opening induces an inhibition of the respiration rate [12]. We show here that these two important features are sufficient to account for the various mCICR profiles, i.e. for the generation of  $\text{Ca}^{2+}$  oscillations and spikes by mitochondria. Our model does not take into account the  $2\text{Na}^+/\text{Ca}^{2+}$  or  $\text{Ca}^{2+}/2\text{H}^+$  exchangers which are not operative in our experimental generation of calcium spikes [9,10]. Although calcium oscillations are rarely observed in mitochondria and not in the conditions in which we observe calcium spiking, our aim in simulating calcium oscillations is to show what parameters have to be varied in order to obtain an excitable organelle exhibiting calcium spiking behaviour.

The simulation of the mitochondrial response to  $\text{Ca}^{2+}$  stimulation corresponding to the parameter values given above, is presented in Fig. 2. After  $\text{Ca}^{2+}$  addition,  $\text{Ca}^{2+}$  is taken up by mitochondria according to Eq. (2), and induces the following processes: the electrogenic uptake of  $\text{Ca}^{2+}$  decreases  $\Delta\Psi$ , this decrease causes an increase of the respiration rate, according to Eq. (7) and Eq. (8), and thus, a  $\text{H}^+$  efflux leading to matrix alcalinization. Such increase in  $\text{pH}_i$

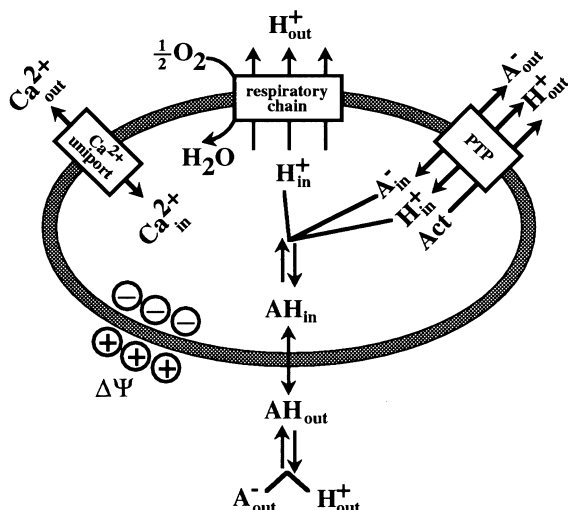


Fig. 1. Schematic representation of the model components. See model description in the text.

triggers PTP opening according to the function  $f$  and Eq. (3) and Eq. (4). Prior to PTP opening, the activation of respiration leads to a slight transient increase of  $\Delta\Psi$ , but then, PTP opens and anion and  $H^+$  fluxes proceed (Eq. (5) and Eq. (6)), resulting in a decrease of  $\Delta\Psi$ , and thus, in the efflux of  $Ca^{2+}$  through the uniporter.  $Ca^{2+}$  efflux proceeds until the equilibrium state between internal exchangeable  $Ca^{2+}$  and external  $Ca^{2+}$  concentration is reached. Simultaneously,  $pH_i$  partially recovers due to the decrease of respiration rate caused by the efflux of activator, and to the influx of  $H^+$  through the PTP. However, in this simulation,  $pH_i$  recovery is not sufficient to close the PTP. Fig. 2 thus simulates the PTP sustained opening mode [15].

$Ca^{2+}$  oscillations and spikes due to mCICR depend on a PTP transitory opening mode [9,12,13], that our model does not simulate with the parameter values used in Fig. 2. Under the conditions of Fig. 2, a very weak efflux of respiration activator proceeds during PTP opening, thus a high respiration rate and a high  $H^+$  efflux are maintained during PTP opening, that sustain a high  $pH_i$ . This is the reason why the PTP remains permanently open. Paradoxically, the only way to obtain a re-uptake of  $Ca^{2+}$  is to slow down the respiratory rate to decrease  $pH_i$ .

This can be achieved by increasing the efflux of the respiration activator through the PTP during channel opening. (Fig. 3). As a result of the faster activator loss during PTP opening, respiration remains inhib-

ited for a longer time, which allows a significant decrease of  $pH_i$ , and thus, the partial closure of the PTP. PTP closure then stops the activator efflux, and then, according to our hypothesis of continuous activator supply, the respiratory rate starts increasing again, as well as  $pH_i$ , which results in PTP re-opening. Then the cycle (Fig. 4) repeats, and leads to  $Ca^{2+}$  oscillations.  $Ca^{2+}$  initiates the oscillatory process by acceleration of respiration and by converting the electric part of  $\Delta\Psi$  into an increase in  $pH_i$  that leads to PTP opening. This is the way by which the activation status of the PTP is connected with the respiratory rate. The transition from monotonous kinetics of  $Ca^{2+}$  efflux (Fig. 2) to damped oscillations (Fig. 3) as observed in [12] is obtained by increasing the respiration inhibition during PTP opening, thus decreasing  $pH_i$  to values that allow the partial closure

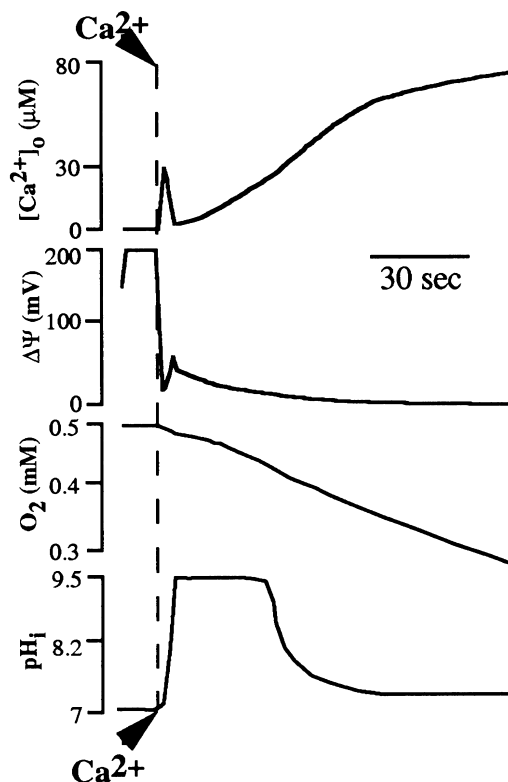


Fig. 2. Simulation of sustained mCICR in response to an external  $Ca^{2+}$  pulse. Extramitochondrial  $Ca^{2+}$  concentration at the moment of stimulation:  $30 \mu M$ . All parameter values are presented in the text (model description). Kinetics of extramitochondrial  $Ca^{2+}$  concentration ( $[Ca^{2+}]_0$ ),  $\Delta\Psi$ , extramitochondrial oxygen ( $O_2$ ), and  $pH_i$ . Concentration in protein is set equal to  $2 \text{ mg ml}^{-1}$ .

of the PTP. The damping can be decreased by globally increasing the membrane and channel permeabilities, leading to sustained oscillations (Figs. 5 and 6) as experimentally observed in [13].

We will now consider the mCICR triggering features exhibited by the model. Fig. 6 shows a simulation of the mitochondrial response to  $\text{Ca}^{2+}$  stimulation pulses of various amplitudes, using the same settings as in Fig. 5. If the amount of added  $\text{Ca}^{2+}$  is sufficient for inducing PTP opening (Fig. 6A), the mitochondrial response varies only slightly when the amount of  $\text{Ca}^{2+}$  used as stimulus is further increased. For instance, the trace shown in Fig. 5 was obtained using twice more  $\text{Ca}^{2+}$  as stimulus than in Fig. 6A (respectively, 30 and 16  $\mu\text{M}$   $\text{Ca}^{2+}$ ). On the contrary, a very slight decrease of the stimulus under a threshold value (15.5  $\mu\text{M}$   $\text{Ca}^{2+}$  in Fig. 6B, instead of 16  $\mu\text{M}$   $\text{Ca}^{2+}$  in Fig. 6A), results in quite different response

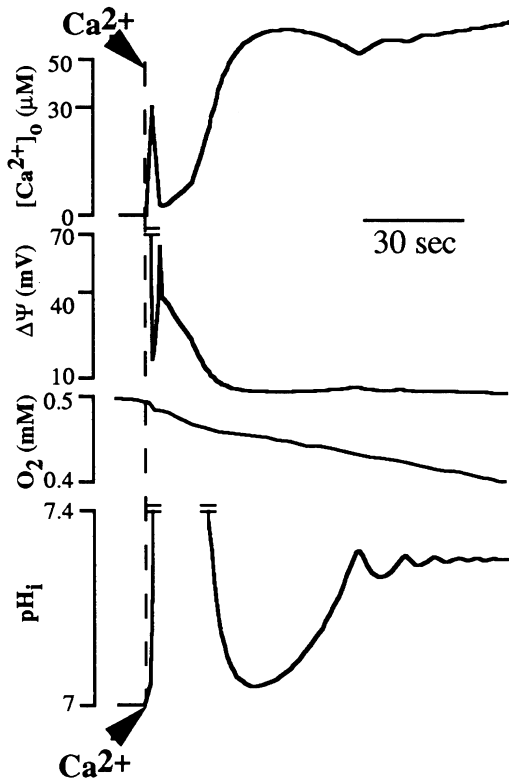


Fig. 3. Simulation of a damped mCICR oscillatory regime, obtained by increasing the efflux of respiration activator during PTP opening. Variables and parameter values are as in Fig. 2, except  $P_A = 20 \text{ min}^{-1}$ .

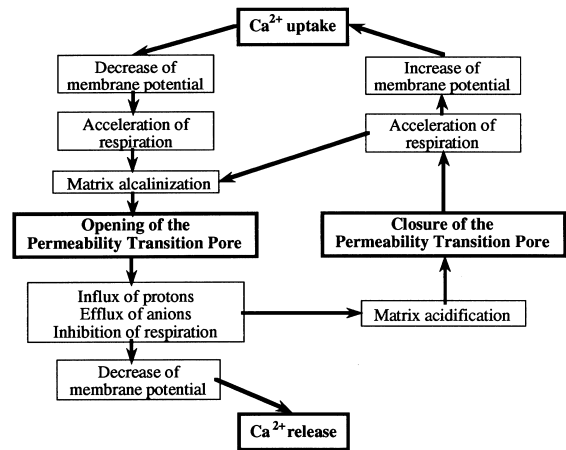


Fig. 4. Sequence of events leading to mCICR.

(Fig. 6B). In the latter case, the stimulus does not trigger PTP opening, and the  $\text{Ca}^{2+}$  provided is simply buffered by mitochondrial  $\text{Ca}^{2+}$  uptake. The model thus simulates the two main features of mCICR, i.e. the threshold-dependence of mCICR triggering, and the all-or-nothing nature of the mCICR response [9].

Fig. 7 shows that the model can also simulate the dependence of mCICR triggering on the frequency of  $\text{Ca}^{2+}$  stimulation pulses (see Fig. e,f in Ref. [9]).

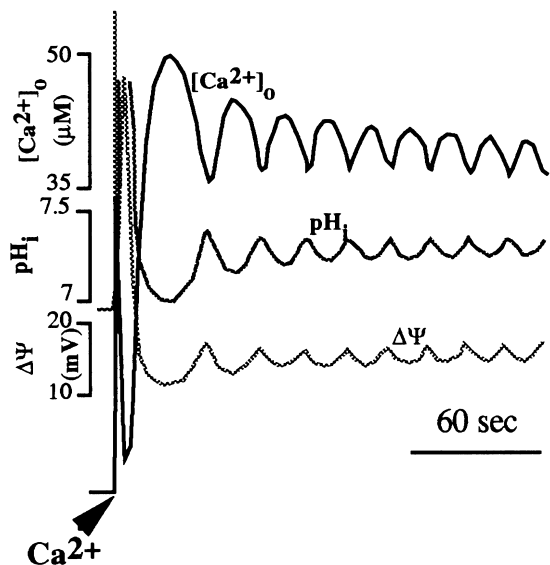


Fig. 5. Simulation of a sustained mCICR oscillatory regime. Parameter values are as in Fig. 3, except  $P_A = 500 \text{ min}^{-1}$ ,  $P_H = 500 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ ,  $P_{AB} = 0.006 \text{ min}^{-1} \text{ mV}^{-1} \text{ mg protein}^{-1}$ ,  $k_{AH} = 300 \text{ min}^{-1} \text{ mg protein}^{-1}$ ,  $k_{f8}^{sp} 8 \text{ min}^{-1}$ .



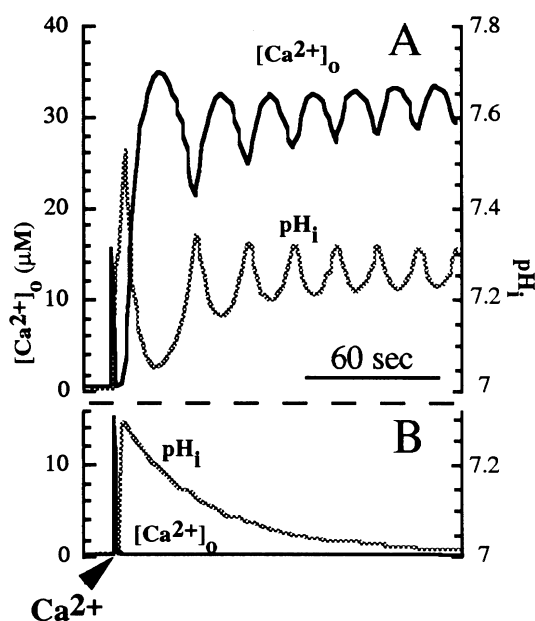


Fig. 6. Simulation of  $[Ca^{2+}]_o$  and  $pH_i$  kinetics in response to different amplitudes of  $Ca^{2+}$  stimulation. Parameter values are as in Fig. 5. (A)  $Ca^{2+}$  stimulation  $16 \mu M$ , (B)  $Ca^{2+}$  stimulation  $15.5 \mu M$ .

Experimentally, in response to a fixed cumulated amount of  $Ca^{2+}$  delivered as a series of pulses, mitochondria undergo mCICR depending on the frequency of the stimulation pulses. This means that the trigger-

ing threshold of mCICR is not a constant in terms of  $Ca^{2+}$  amount, but depends on the rate at which  $Ca^{2+}$  is provided to mitochondria [9]. Fig. 7 shows a simulation of this feature. Three stimulations with a period of 0.5 min do not induce PTP opening (Fig. 7B), while the same stimulations performed with a period of 0.2 min induce PTP opening (Fig. 7A), and further  $Ca^{2+}$  oscillations. The amount of  $Ca^{2+}$  taken up by mitochondria at the end of the stimulation sequence is the same in both cases, as well as the value of  $\Delta\Psi$ . In both cases, the amplitude of the  $\Delta\Psi$  drops, that respond to the stimulation pulses, is constant. A very similar situation is observed experimentally [9]. On the other hand, the recovery of  $pH_i$  is a much slower process than recovery of  $\Delta\Psi$  or than the uptake of added  $Ca^{2+}$ . Therefore, the higher the frequency of stimulation, the higher the initial  $pH_i$  value at the time of each stimulation pulse. Thus, among the factors triggering PTP opening,  $pH_i$  is sufficient to account for the rate-dependence of mCICR triggering. Our recent experimental observations [10], and the present model analysis, show that  $pH_i$  is the key factor regulating the PTP activation status during mCICR.

As shown in Fig. 6, the probability of PTP opening sharply depends on the initial  $Ca^{2+}$  amount used as stimulus, because the latter is connected stoichiometrically with  $H^+$  respiratory efflux. However, the maximal  $pH_i$  level reached during  $Ca^{2+}$  uptake depends

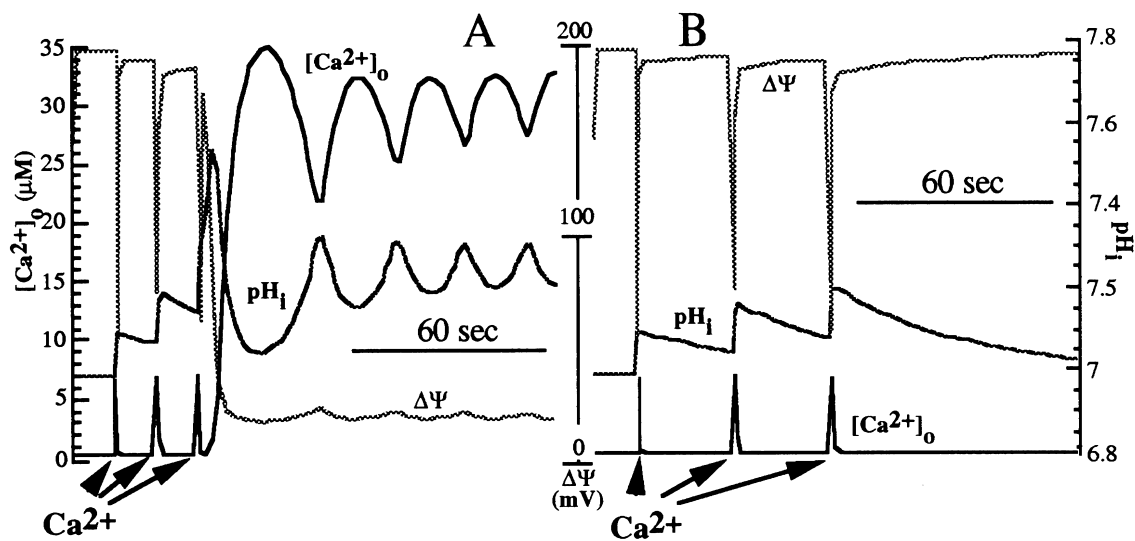


Fig. 7. Simulation of  $[Ca^{2+}]_o$ ,  $\Delta\Psi$ , and  $pH_i$  kinetics in response to series of under-threshold  $Ca^{2+}$  pulses of variable frequencies. Value of each  $Ca^{2+}$  pulse is  $7 \mu M$ . (A) Time interval between pulses 0.2 min. (B) Time interval between pulses 0.5 min; other parameters as in Fig. 5.

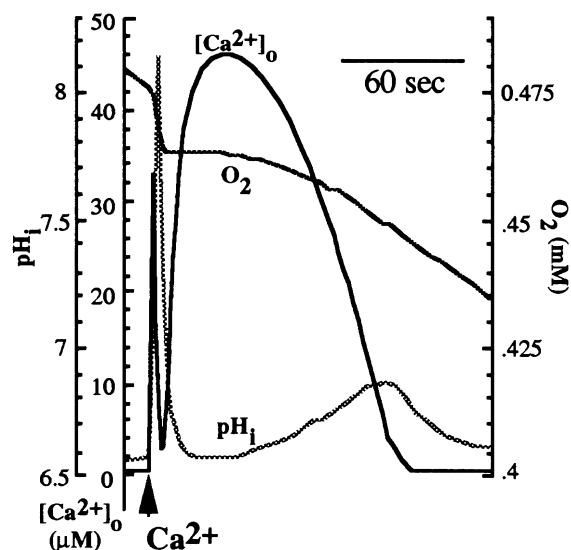


Fig. 8. Simulation of a mCICR spike, obtained by increasing the electroneutral flux of weak acids;  $k_{AH} = 300 \text{ min}^{-1} \text{ mg protein}^{-1}$ . Other parameters as in Fig. 6A.

also on the opposite electroneutral flux of weak acids (Eq. (13)) which buffers the  $\text{pH}_i$  changes. This effect is particularly apparent in Fig. 7: due to the electroneutral flux of weak acids,  $\text{pH}_i$  recovers when the time interval between stimulations is long, thus, mitochondria can take up more  $\text{Ca}^{2+}$  without activation of the PTP. The decrease of the respiration rate that occurs after PTP opening also contributes to  $\text{pH}_i$  recovery. Indeed, the slower the respiration rate, and hence the smaller its ratio to the electroneutral flux, the more  $\text{Ca}^{2+}$  taken-up before  $\text{pH}_i$  reaches the value triggering PTP opening. At slower respiratory  $\text{H}^+$  efflux and faster opposite electroneutral flux, the equilibrium between this two fluxes can be reached at low  $\text{pH}_i$  values that do not trigger PTP opening. One can suppose that in these conditions, all the  $\text{Ca}^{2+}$  released could then be taken-up without re-activation of the PTP. Fig. 8 presents such a simulation, where the respiration rate is low and the electroneutral flux of weak acids high. In these conditions, a single  $\text{Ca}^{2+}$  spike is simulated. Immediately after  $\text{Ca}^{2+}$  stimulation, the respiration rate strongly increases, resulting in a consequent increase of  $\text{pH}_i$ . Then, after  $\text{Ca}^{2+}$  release, only a slight increase of the respiration rate and of  $\text{pH}_i$  takes place, although all released  $\text{Ca}^{2+}$  is taken up again. Such single spikes represent an important mode of mCICR [9].

Further improvements of the model are currently studied in our laboratory: they concern (i) the modelling of the calcium fluxes through the calcium uniporter; recently Magnus and Keizer [27] introduced an allosteric term in the calcium flux equation in order to account for data of Wingrove et al. [20]. This is certainly a better model to account for the threshold observed in the calcium concentration for calcium influx.

(ii) Another important point concerns the modelling of the calcium pools inside mitochondria; a  $100 \mu\text{M}$  free internal calcium is certainly unrealistic, but however the amount of calcium uptake by mitochondria is huge under bound form; thus a bound calcium pool has to be considered that becomes releasable after a  $\Delta\Psi$  collapse.

The model presented here however, simulates all the mCICR profiles observed experimentally, as well as the notable triggering characteristics of mCICR. Since mCICR contributes to cell  $\text{Ca}^{2+}$  signalling [9,10], the model we present here should appear useful to further mathematically address the role of mCICR in the spatiotemporal organisation of intracellular  $\text{Ca}^{2+}$  signals as already proposed by Meyer and Stryer [28].

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